INFLUENCE OF SHOCK WAVES ON THE LIGHT CURVES OF LONG-PERIOD VARIABLES

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ABSTRACT. The shapes of light curves of Mira-type variable stars is analyzed. It is shown that some features of the light curves may be explained by passage of shock waves in the stars' atmospheres. In particular, it is noted that "steps" of the ascending branches of the light curves have very similar durations (of the order of 0.1 stellar period) for a big number of miras. This time span is just approximately the time of propagation of shock waves through stars' atmospheres. The numerical simulation of shock wave propagation support the above suggestion and yield time scales that fit the observed light curve features.

Key words: Stars: Miras, light curves, stellar atmospheres, shock waves.

For long-period variables, the shape of light curves is determined to a considerable extent by influence of shock waves, which propagate in the star's atmosphere and in the inner layers of the circumstellar envelope during each cycle of the stellar variability. Of interest are also some peculiarities of the light curves: "humps", "steps" on the ascending branch and - as an extreme case - double maxima. can be said that the "normal", stable state of a long-period variable is the minimum light, whereas the shock waves, caused by stellar pulsations, move the star from this state to the "excited" state, which is connected with the light maximum. The increase of brightness in the visible and infrared is in the first turn due

ABSTRACT. The shapes of light curves of to an increase of temperature of the photosphelira-type variable stars is analyzed. It is ric layers that are responsible for the stellar from that some features of the light curves continuum in the visual.

> In order to investigate the influence of shock waves on the light curves, we constructed a numerical model of propagation of a shock in mira's atmosphere. We accounted for the effects of ionization of gas and for dissociation of molecular hydrogen on the shock front. We assumed the shock to be spherical. The initial radius was $\sim 3 \cdot 10^{13}$ cm. Then, with a fixed step in time (10^5-10^6 s) , we computed the gas physical parameters for the current shock radius. According to Willson (1976), we assumed that the shock velocity falls with growing radius as $D = D_0(r_0/r)^{\alpha}$; from empirical data, $\alpha = 1$. Here D is the current shock velocity, D_0 - its initial value at $r = r_0$. We varied D_0 between 20 and 100 km s⁻¹. We assumed also the postshock ionization and dissociation to be determined by Saha-type equations.

We solved by the iteration method the system of equations for the quantities

$$z = \frac{p_2}{p_1} = \frac{2\gamma_1 \mathbf{M}^2}{\gamma_2 + 1} \left[1 + \frac{(\gamma_2^2 - 1) q}{2D^2} \right],$$

$$\frac{T_2}{T_1} = \frac{2\gamma_1 (\gamma_2 - 1) \mu_2 \mathbf{M}^2}{(\gamma_2 + 1)^2 \mu_1} \times \left[1 - \frac{(3 - \gamma_2) (\gamma_2 + 1) q}{2D^2} \right],$$
(1)

where q is the energy spent for ionization and dissociation, M is the Mach number, γ - the adiabatic index, μ - the molecular we-

Table 1: Characteristics of individual stars. Notes: Gr. – group of periods $(I - P < 200^d; II - 200^d < P < 300^d; III - 300^d < P < 400^d; IV - P > 400^d); f$ – asymmetry of the light curve; T – duration of the hump; ' – the period was redetermined.

| as redetermi Star | P, days | Gr. | Sp. (latest) | f | lg(T/P) | T/P | $\lg[\Delta P/P]$ |
|----------------------|---------|-----|--------------|------|--------------------------|----------------------|-------------------|
| R Hor | 407.6 | IV | M8eII-III | 0.40 | $\frac{16(1/1)}{-0.844}$ | $\frac{1/1}{0.1431}$ | ig[ΔF/F] |
| U Aur | 408.09 | IV | M9e | 0.39 | -0.845 | 0.1431 | |
| S Scl | 362.57 | III | M9e (Tc) | 0.48 | -1.002 | 0.0996 | |
| V Cyg | 421.27 | IV | C7.4e | 0.46 | -1.002 | 0.0989 | |
| S Cep | 502.34 | ΙV | C7.4e | 0.55 | -0.864 | 0.1369 | -1.498' |
| T Cep | 388.14 | III | M8.8e | 0.54 | -0.915 | 0.1303 0.1217 | -1.430 |
| R Cas | | IV | M10e | 0.40 | -1.111 | | -1.506' |
| W Peg | 345.5 | Ш | M8e | 0.46 | -0.919 | 0.1206 | -1.500 |
| RU Tau | 544.6 | IV | M6.5 | 0.62 | -0.845 | 0.1200 0.1428 | |
| V Cam | 503.16 | ĪV | M7e | 0.31 | -0.973 | 0.1428 | -1.433 |
| χ Cyg | 421.54 | ĪV | S10.4e(MS) | 0.41 | -1.167 | 0.0681 | -1.433 |
| R Aur | 448.10 | īV | M9.5e | 0.51 | -0.874 | 0.0031 0.1336 | -1.686 |
| Y Cep | 322.57 | III | M8.2e | 0.40 | -1.124 | 0.0752 | -1.000 |
| W And | 399.24 | III | M10(S9.2e) | 0.42 | -0.923 | 0.0132 | -2.076' |
| o Cet | 333.47 | III | M9e | 0.38 | -1.036 | 0.0921 | -2.347' |
| RR Cep | 384.18 | Ш | M8.8e | 0.41 | -0.862 | 0.0321 0.1374 | -2.047 |
| RT Oph | 426.34 | ΙV | M7e(C) | 0.36 | -1.107 | 0.1314 | |
| U Her | 415.69 | IV | M9.5e | 0.40 | -1.019 | 0.0152 | -1.627' |
| NP Her | 448 | ĪV | C6.3e | 0.5 | -1.013 | 0.0992 | -1.024 |
| RU Her | 484.83 | ĪV | M9 | 0.43 | -0.941 | 0.0332 | |
| U Ser | 237.50 | II | M6e | 0.48 | -0.591 | 0.2563 | |
| W Cas | 417.98 | ΙV | C7.1e | 0.46 | -0.891 | 0.2365 0.1286 | -1.444' |
| T Phe | 281.79 | 11 | M5e | 0.37 | -0.973 | 0.1265 | 1.111 |
| R Cam | 270.22 | II | S8.7e | 0.45 | -0.727 | 0.1876 | |
| RS Her | 219.70 | II | M8: | 0.47 | -0.919 | 0.1206 | |
| VZ Cas | 169.24 | I | МЗе | 0.46 | -0.985 | 0.1034 | |
| RS UMa | 258.97 | II | M6e | 0.42 | -1.178 | 0.0664 | |
| R Vir | 145.95 | I | M8.5e(III) | 0.50 | -1.204 | 0.0625 | -2.678' |
| TU And | 318.36 | III | M5e | 0.48 | -1.140 | 0.0725 | -2.301' |
| R Leo | 313.67 | III | M9.5eIII | 0.43 | -1.307 | 0.0493 | -1.921' |
| U CMi | 422.03 | IV | M4e | 0.52 | -0.768 | 0.1707 | -1.706' |
| X Cas | 441.94 | IV | C5.4e | 0.55 | -0.658 | 0.2200 | -1.345' |
| T Cas | 465.35 | IV | M9.0e | 0.56 | -0.894 | 0.1275 | -1.336' |
| R UMa | 301.53 | III | M9e | 0.39 | -1.342 | 0.0455 | -3.523' |
| T Cam | 374.80 | III | S8.5e | 0.47 | -1.135 | 0.0732 | -3.323 $-2.367'$ |
| T UMi | 316.74 | III | M6e | 0.45 | -1.313 | 0.0486 | -2.301 $-1.281'$ |
| S Umi | 225.87 | II | S5.9e | 0.47 | 1.010 | 0.2117 | -1.201 |
| R Aqr | 386.96 | III | M8.5e | 0.42 | | 0.1236 | |
| X UMa | 249.04 | II | M4e | 0.43 | | 0.1230 0.0524 | |
| S Leo | 190.16 | I | M6e: | 0.47 | | 0.0324 0.1829 | |
| T Lyn | 406.0 | IV | C7.1e(N0e) | 0.47 | | 0.1071 | |
| U Lyn | 433.6 | IV | M9.5:e | 0.42 | | 0.1604 | |
| SZ Aur | 454.04 | ĪV | M8e | 0.46 | | 0.1004 0.2394 | |
| | 101.01 | 4 7 | 1.100 | 0.40 | | 0.2394 | |

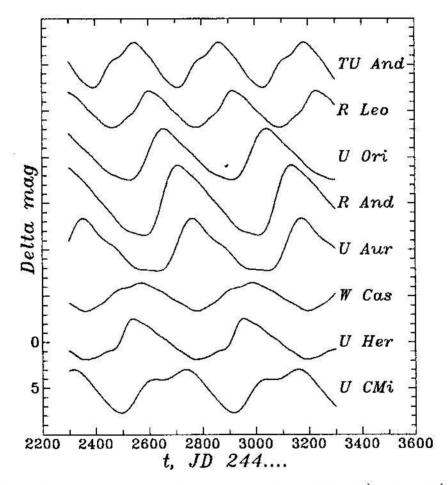


Figure 1: The smoothed light curves of some Mira-type stars with periods $300^d < P < 400^d$. The observations were obtained by the AAVSO members (Mattei et al. 1978). For approximation we have used a program FOUR-N by Andronov (1994) which determines a least squares fit with different number of harmonics.

| Table 1 (continued) | | | | | | | | | | | |
|---------------------|---------|-----|--------------|----------------|-------------|--------|-------------------|--|--|--|--|
| Star | P, days | Gr. | Sp. (latest) | \overline{f} | $\lg (T/P)$ | T/P | $\lg[\Delta P/P]$ | | | | |
| V Del | 533.51 | IV | M6e | 0.42 | 80000 00 | 0.0978 | | | | | |
| U Ari | 371.13 | III | M9.5e | 0.40 | | 0.1172 | | | | | |
| S CMi | 332.94 | III | M8e | 0.49 | | 0.1045 | | | | | |
| R CVn | 328.53 | III | M9e | 0.46 | | 0.1191 | | | | | |
| R LMi | 372.19 | Ш | M9.0e(Tc:) | 0.41 | | 0.0935 | | | | | |

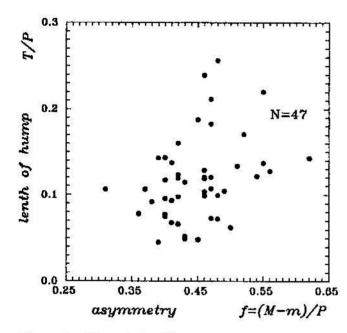


Figure 2: The relationship "duration of the hump -asymmetry". NP Her is not included.

ight. Subscripts 1 and 2 refer to the pre—and postshock gas parameters respectively. Detailed expressions for the quantities, entering the above formulae, can be found in Klimishin (1984). The calculations stopped when the postshock degree of ionization fell to a value $a = N_{\rm H^+}/N_{\rm H} < 0.001$. As a byproduct, we got also the intensity of the free–free emission of the postshock ionized gas.

The main result of the calculations is the time of the shock propagation in the stellar atmosphere until the moment when the shock loses its capacity to heat the gas to temperature that is sufficient for ionization. This time span is on the average about 1 month (from 20 to 35 days). This value is close to the duration of "humps" and "steps" on the ascending branches of miras' light curves. The subsequent growth of brightness is going on as if by inertia, at the expense of the growing volume of the expanding heated gas. It should be stressed that the relative duration of the noted features, expressed in fractions of the light period, is notably constant for different stars: it is, on the average, ~ 0.1 .

The figures present sample light curves and some statistical regularities, which imply the plausibility of the suggested interpretation of the light curves. The table lists the stars used in the statistics.

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